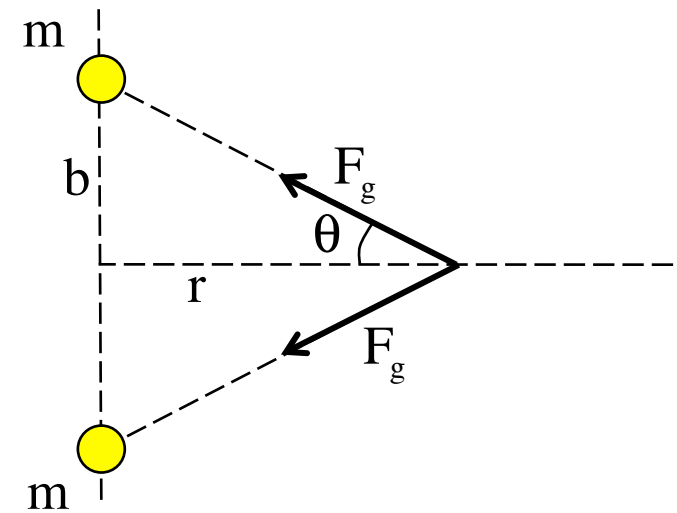


## Problem 13.26

Determine the gravitational field “ $r$ ” units down the  $x$ -axis due to the two masses shown.

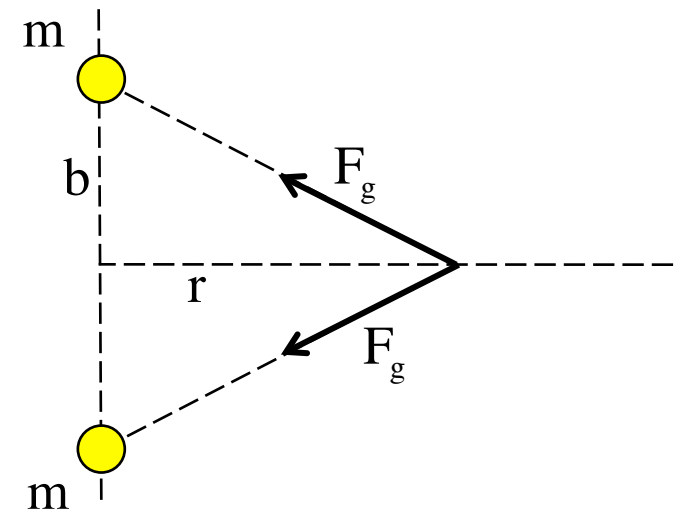
Because the masses are the same, and due to the symmetry of the situation, the magnitude and angle of the accelerations acting at “ $r$ ” will be the same. As such, the  $y$ -components will add to zero and all we will have to calculate is the  $x$ -component of one, then double to accommodate the second. That is:



$$\begin{aligned} a_x &= 2 \left( \frac{Gm}{\left( (b^2 + r^2)^{1/2} \right)^2} \right) \cos \theta \\ &= 2 \left( \frac{Gm}{\left( (b^2 + r^2)^{1/2} \right)^2} \right) \left( \frac{r}{(b^2 + r^2)^{1/2}} \right) \\ &= 2 \frac{Gm}{\left( (b^2 + r^2)^{3/2} \right)} r \end{aligned}$$

b.) Why should the net field go to zero as “r” goes to zero?

As “r” goes to zero, the *x*-components of the two forces becomes less and less until they are zero, and the *y*-components always add to zero as they are equal and opposite in direction.



c.) Prove *Part b*'s answer mathematically.

We derived the acceleration in the *x*-direction, as shown below.

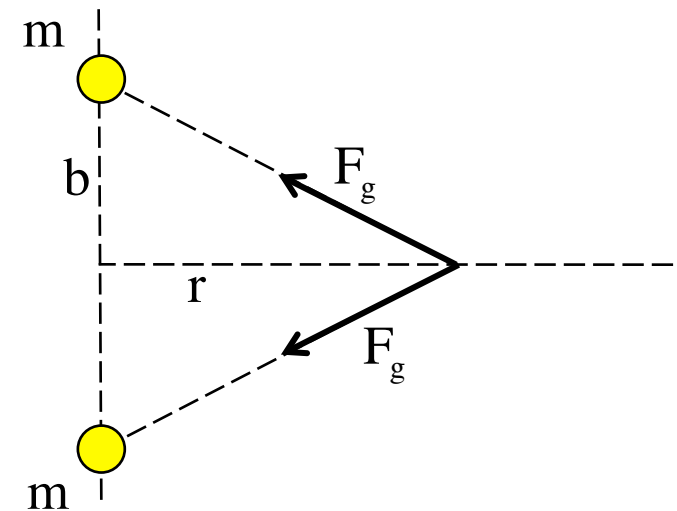
$$a_x = 2 \left( \frac{Gm}{\left( (b^2 + r^2)^{1/2} \right)^2} \right) \cos \theta$$

As  $r \Rightarrow 0$ ,  $\theta \Rightarrow 90^\circ$  and  $\cos 90^\circ = 0$ .

d.) If we let “r” go to infinity, what happens?

As you get farther and farther away, the system begins to resemble more and more that of a point mass of mass “2m.” With that, the field should be:

$$a_x = \frac{G(2m)}{r^2}$$



e.) prove your response to *Part d.*

If you note that as “r” gets really big, the contribution of “b” in the expression gets less and less significant and essentially becomes nil, and we can write:

$$\begin{aligned} a_x &= 2 \frac{Gm}{(b^2 + r^2)^{3/2}} r \\ &= 2 \frac{Gm}{(r^2)^{3/2}} r \\ &= 2 \frac{Gm}{r^2} \end{aligned}$$